

Theoretical Studies on Dimerization Reactions of 4, 7-diphenyl-1,10-phenanthroline (BPhen) and Bathocuproine (BCP) in Organic Semiconductors¹

Fahimeh Shojaie

Department of Photonics, Graduate University of Advanced Technology, Kerman, Iran

e-mail: fahimeh_shojaie@yahoo.com

Submitted September 11, 2013; accepted for publication, November 5, 2013

Abstract—The dimerization reactions of 4, 7-diphenyl-1,10-phenanthroline (BPhen) and bathocuproine (BCP) were studied theoretically. In this work, the molecular geometries, BPhen and BCP monomers, positively charged BPhen and BCP intermediates, and BPhen and BCP dimers were calculated in neutral state at different levels of ab initio methods for the dimerization reactions. The molecular geometries of the BPhen and the BCP monomers and their dimers in their cationic and anionic states were optimized using the B3LYP/6-31G(d) method. The π orbitals interactions in the phenyl groups are responsible for the shortening of the C₆–N₁ and C₇–N₁ bond distances in BPhen and BCP dimers. The large difference between the reorganization energy components of BCP dimer for hole transport process is due to larger energy required to reorganize C₅C₄N₂C₆ and C₈C₁N₁C₇ angles. The methyl groups and the phenyl rings act as sterically hindering groups in the BPhen and BCP dimers. In contrast to BPhen dimer where its electron hopping rate is less than BPhen monomer, the magnitude of electron hopping rate of BCP dimer is 12 times higher than BCP monomer.

DOI: 10.1134/S1063782614080223

1. INTRODUCTION

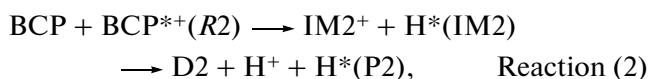
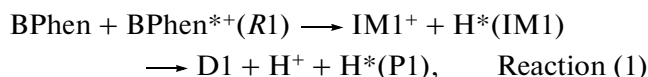
It has been demonstrated that organic light-emitting diodes (OLEDs) can be used for flat panel displays and solid-state lighting applications [1, 2]. Intensive researches have been carried out to find materials with high light-emitting efficiencies, high thermal stabilities and good amorphous film formation properties [3, 4]. The optoelectronic properties of OLED devices depend on the highest occupied molecular orbital (HOMO) and the lowest unoccupied molecular orbital (LUMO) energy levels and electron and hole mobility [3, 4].

Sebastian et al. studied photochemical reactions of different organic semiconductor materials commonly used in OLEDs by means of laser desorption/ionization time-of-flight mass spectrometry (LDI–TOF–MS) technique [5]. These photochemical reactions were induced in the organic semiconductor materials by UV-laser light irradiation under the influence of a positive bias voltage [5]. The analysis of LDI–TOF–MS spectra allows an additional understanding of the underlying degradation mechanisms, which may occur in electrically driven organic semiconductor devices [5]. Sebastian et al. studied possible reaction pathways and the different stages of dimerization of the well-known electron transport material BPhen

and compared these to the results with its analog BCP, showing the degradation of the two materials under hole transporting conditions [5]. Several studies have been carried out over the past years revealing different degradation mechanisms for OLEDs [6–13].

Sebastian et al. applied laser desorption/ionization (LDI) technique to create the dimerization reactions of the BPhen and BCP monomers and showed these reactions occurred via the pathway that is shown in Fig. 1 [5].

The pathway for the dimerization reactions in both cases (BPhen and BCP) consists of two reaction steps. The first step is a neutral monomer (BPhen, BCP) addition to an excited positively charged monomer molecule (BPhen^{*+}, BCP^{*+}) and the formation of the intermediates (IM1, IM2). The second step is the production of products (P1, P2). The reaction pathways in Fig. 1 can be written in short forms as follows:



where the D1 and the D2 are the BPhen and the BCP dimers respectively.

The thermodynamic properties of these reaction pathways were calculated in the standard states using a

¹ The article is published in the original.

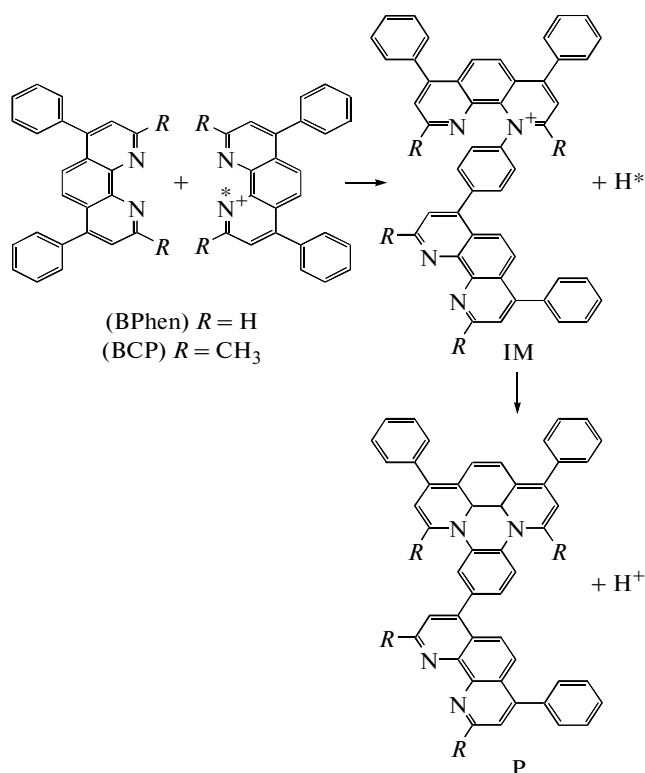


Fig. 1. Dimerization reactions of BPhen ($R = H$) and BCP ($R = CH_3$) after a short laser pulse excitation and an applied positive high voltage by LDI experiments [5].

combination of quantum mechanical computations performed by the Gaussian 03 program [14] at the HF/3-21G*, HF/6-31G(*d*), B3LYP/3-21G* and B3LYP/6-31G(*d*) levels. Molecular orbital (MO) calculations were performed at the same levels too.

The molecular geometries of the monomers (BPhen, BCP) and their dimers (D1 and D2) in their cationic and anionic states were optimized by the Density functional theory (DFT) modelling method, DFT/B3LYP, using the 6-31G(*d*) basis set. The ionization potentials, electron affinities and reorganization energies of the above compounds were calculated using B3LYP/6-31G(*d*) method. In this paper, we present DFT calculations to determine the hopping transport rate of electron and hole. The charge carrier transport is an important factor for developing high efficiency OLED materials.

2. COMPUTATIONAL METHODS

The molecular geometries of the monomers (BPhen, BCP), positively charged intermediates (IM1⁺, IM2⁺) and the dimers (D1, D2) were optimized at the HF/3-21G*, HF/6-31G(*d*), B3LYP/3-21G* and B3LYP/6-31G(*d*) levels of ab initio method.

Vibrational frequencies of these compounds, also calculated at the same levels of ab initio method, have

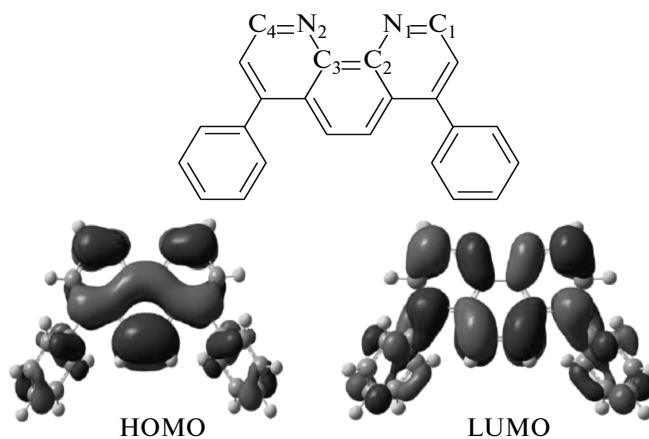


Fig. 2. The structure of the BPhen in neutral state and the electron density isocontours of HOMO and LUMO using the DFT/B3LYP method with the 6-31G(*d*) basis set.

been used to determine zero-point energy (ZPE), entropy, free energy and enthalpy. The energy of the compounds excited state was calculated by the time-dependent self-consistent-field (TD-SCF) method [15]. The ionization potentials, electron affinities and reorganization energies of these compounds have been calculated at B3LYP/6-31G(*d*) level. All calculations were carried out using the Gaussian 03 program.

3. RESULTS AND DISCUSSION

3.1. Optimized Molecular Geometries

The structures and the electron density isocontours of HOMO and LUMO of the BPhen and BCP monomers, positively charged intermediates IM1 and IM2 and the dimers D1 and D2 are shown in Fig. 2–7. The optimized molecular geometries, the energies of frontier molecular orbitals (E_{LUMO} , E_{HOMO}) and the HOMO and LUMO energy gaps ($\Delta E = E_{LUMO} - E_{HOMO}$) these structures are shown in Tables 1–6.

Tables 1–6 show that the molecular geometries that were calculated with different basis sets are only slightly different. The electron density isocontours (Figs. 3 and 6) show that in both IM1⁺ and IM2⁺ intermediates, the π orbitals of the star labeled rings make a large contribution to their LUMO energies whereas these π orbitals do not contribute to their HOMO energies. In case of the D1 and D2 dimers, the density isocontours show that the π orbitals of the star labeled rings make major contributions to both LUMO and HOMO energies of the dimers (Tables 3 and 6). When an H⁺ is removed from IM1⁺ and IM2⁺, C₆–N₁ and C₇–N₁ bond distances are shortened by 1.787 and 1.692 Å, respectively. The π orbitals interactions in the phenyl groups are responsible for the shortening of the C₆–N₁ and C₇–N₁ bond distances. Since one of the most accurate theoretical techniques is the density functional theory which can be applied to most sys-

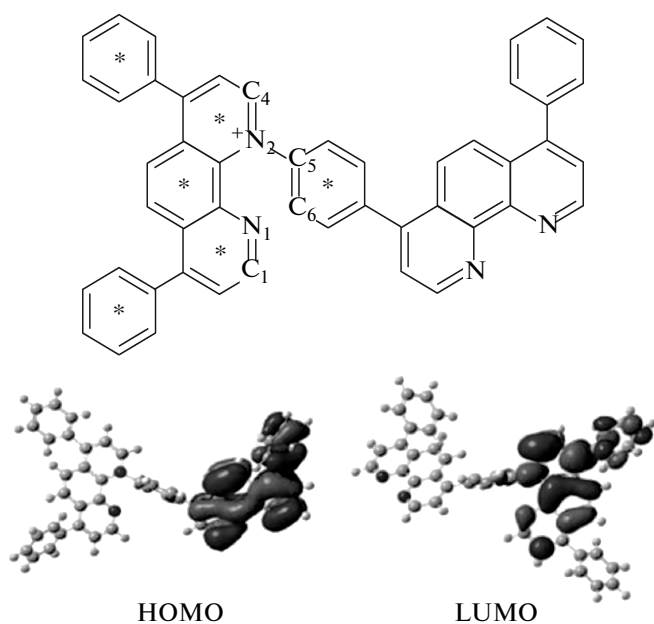


Fig. 3. The structure of the IM1⁺ in neutral state and the electron density isocontours of HOMO and LUMO using the DFT/B3LYP method with the 6-31G(*d*) basis set.

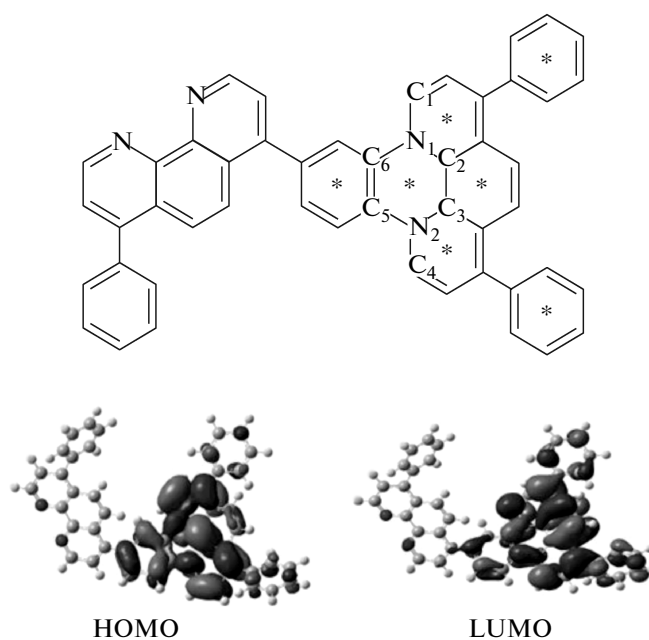


Fig. 4. The structure of the D1 in neutral state and the electron density isocontours of HOMO and LUMO using the DFT/B3LYP method with the 6-31G(*d*) basis set.

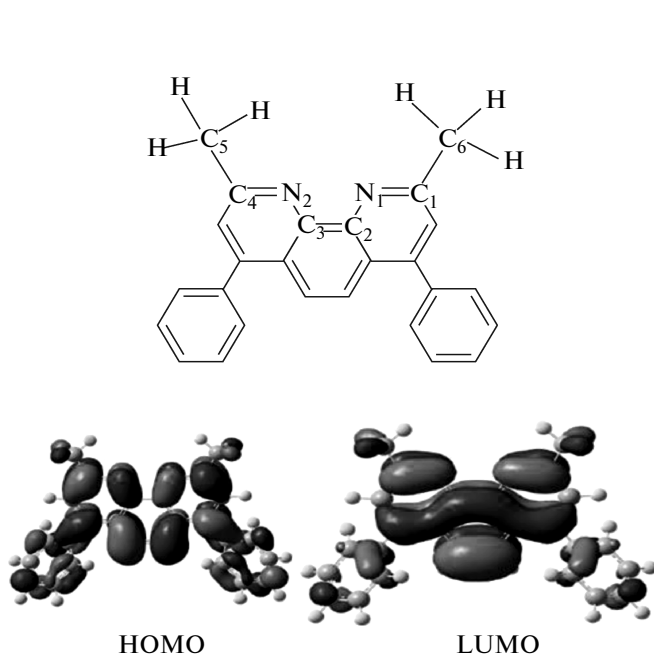


Fig. 5. The structure of the BCP in neutral state and the electron density isocontours of HOMO and LUMO using the DFT/B3LYP method with the 6-31G(*d*) basis set.

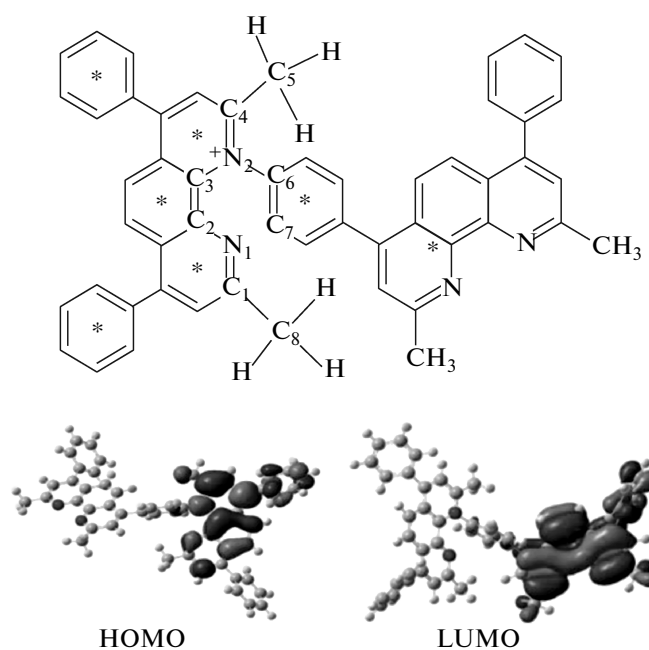


Fig. 6. The structure of the IM2⁺ in neutral state and the electron density isocontours of HOMO and LUMO using the DFT/B3LYP method with the 6-31G(*d*) basis set.

tems and allows the structure to be efficiently optimized [16–18], so B3LYP/6-31G(*d*) clearly provide better results for organic molecules. BPhen has energy gap 3.4 eV [19] and BCP has 4.2 eV [20], as shown in Table 1 and 4, these energies by B3LYP/6-31G(*d*) are in better agreement with these values.

The geometrical variations in the monomers and the dimers can be noticed from their HOMO and LUMO. The dimers have the higher HOMO and the lower LUMO levels than their monomers. The ΔE values of the dimers are lower than their monomers counterparts. The IM1⁺ and the IM2⁺ act as electron traps

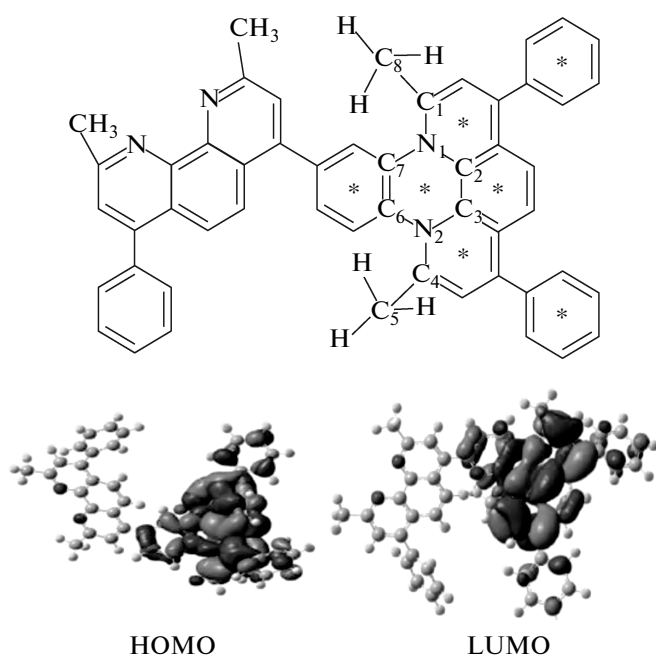


Fig. 7. The structure of the D2 in neutral state and the electron density isocontours of HOMO and LUMO using the DFT/B3LYP method with the 6-31G(*d*) basis set.

because they have the lowest HOMO levels. The HOMO energies of the dimers are larger than the HOMO energies of the monomers. The LUMO energies of the intermediates are smaller than the LUMO energies of both the monomers and the dimers. The geometrical variations in the BPhen and BCP in their cationic and anionic states can be seen by comparing their HOMO and LUMO energies (See Tables 1, 4, 7 and 8).

The HOMO and LUMO of the rings make major contributions to the large geometrical variations in the

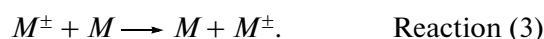
BPhen and BCP. The orbitals of nitrogen atoms contribute to the formation of both the HOMO and LUMO antibondings. The HOMO and LUMO antibondings strongly interact with the π orbitals of the rings. The lengthening of C₁–N₁ and C₄–N₂ and the shortening of C₂–N₁ and C₃–N₂ bond distances in both cationic and anionic states are expected by the molecular orbitals of BPhen and BCP. The presence of the methyl groups in the BCP has negligible effect on the HOMO and LUMO.

Tables D1 (A), D1 (B), D2 (A) and D2 (B) compare some of the bond lengths and torsional angles of the D1 and the D2 dimers in neutral (Tables 3 and 6), cationic and anionic states (Tables 9 and 10).

3.2. Reorganization Energy

The charge transport properties of organic materials are well described in the formalism of Gaussian disorders by Bassler and co-worker [21]. Energy and position disorders are two key parameters in this formalism [22]. For designing suitable hole or electron transporting materials, it is difficult to relate these two disorder parameters to molecular properties [22]. The measured hole and electron mobility of oligoacene crystals show a band hopping transition occurring at about room temperature [23].

Hopping of hole or electron can be described as an electron transfer (ET) from a charged, relaxed unit to an adjacent neutral unit and the mobility depends on the electron transfer rate [23]. The transfer of hole and electron can be described as a self-exchange reaction (3) [23]



Where M^{\pm} refers to the molecule in a cationic or anionic state and M indicates the adjacent molecule in

Table 1. Molecular geometry, bond distances, bond angles, torsional angles, the calculated HOMO and LUMO energies and energy gaps of BPhen in neutral state

Bond/angle name	HF/3-21G*	HF/6-31G(<i>d</i>)	B3LYP/3-21G*	B3LYP/6-31G(<i>d</i>)
C ₂ –C ₃ , Å	1.451	1.461	1.455	1.461
N ₁ –N ₂ , Å	2.744	2.711	2.743	2.717
C ₁ –N ₁ , Å	1.309	1.297	1.332	1.321
C ₂ –N ₁ , Å	1.342	1.339	1.359	1.352
C ₃ –N ₂ , Å	1.342	1.339	1.359	1.352
C ₄ –N ₂ , Å	1.309	1.297	1.332	1.321
N ₁ C ₂ C ₃ N ₂ , deg	0.000	0.000	–0.000	0.000
C ₁ N ₁ C ₂ , deg	119.57	118.38	118.29	117.74
C ₃ N ₂ C ₄ , deg	119.57	118.38	118.29	117.74
E_{LUMO} , eV	2.2653	2.3507	–1.4854	–1.4658
E_{HOMO} , eV	–8.2147	–7.9584	–6.1447	–5.9758
ΔE , eV	10.4800	10.3091	4.6593	4.50995

Table 2. Molecular geometry, bond distances, bond angles, torsional angles, the calculated HOMO and LUMO energies and energy gaps of IM1⁺ in neutral state

Bond/angle name	HF/3-21G*	HF/6-31G(<i>d</i>)	B3LYP/3-21G*	B3LYP/6-31G(<i>d</i>)
C ₂ –C ₃ , Å	1.449	1.457	1.448	1.454
N ₁ –N ₂ , Å	2.839	2.837	2.834	2.837
C ₁ –N ₁ , Å	1.306	1.296	1.331	1.321
C ₂ –N ₁ , Å	1.339	1.337	1.353	1.348
C ₃ –N ₂ , Å	1.372	1.337	1.385	1.388
C ₄ –N ₂ , Å	1.338	1.329	1.359	1.350
C ₅ –N ₂ , Å	1.447	1.465	1.484	1.470
C ₆ –N ₁ , Å	3.159	3.188	3.199	3.210
N ₁ C ₂ C ₃ N ₂ , deg	–2.317	–8.940	–5.179	–9.654
C ₁ N ₁ C ₂ , deg	120.26	118.99	118.93	118.27
C ₃ N ₂ C ₄ , deg	120.88	120.47	120.84	120.61
C ₅ N ₂ C ₃ , deg	125.58	125.98	125.38	125.64
<i>E</i> _{LUMO} , eV	–2.2590	–2.2590	–5.809	–5.6998
<i>E</i> _{HOMO} , eV	–9.6449	–9.6449	–7.7284	–7.6275
ΔE , eV	7.3858	7.3858	1.9194	1.9276

Table 3. Molecular geometry, bond distances, bond angles, torsional angles, the calculated HOMO and LUMO energies and energy gaps of D1 in neutral state

Bond/angle name	HF/3-21G*	HF/6-31G(<i>d</i>)	B3LYP/3-21G*	B3LYP/6-31G(<i>d</i>)
C ₂ –C ₃ , Å	1.328	1.330	1.361	1.363
C ₁ –N ₁ , Å	1.388	1.382	1.396	1.387
C ₂ –N ₁ , Å	1.421	1.416	1.435	1.429
C ₃ –N ₂ , Å	1.421	1.416	1.433	1.427
C ₄ –N ₂ , Å	1.388	1.381	1.393	1.385
C ₅ –N ₂ , Å	1.410	1.408	1.424	1.419
C ₅ –C ₆ , Å	1.407	1.405	1.419	1.413
C ₆ –N ₁ , Å	1.411	1.407	1.418	1.413
N ₁ C ₂ C ₃ N ₂ , deg	0.044	0.023	0.281	–0.017
N ₂ C ₅ C ₆ N ₁ , deg	0.290	0.496	1.366	1.188
C ₂ C ₃ N ₂ C ₅ , deg	20.555	29.268	18.723	26.715
C ₁ N ₁ C ₂ , deg	118.64	118.47	118.47	118.35
C ₃ N ₂ C ₄ , deg	118.65	118.48	118.66	118.48
C ₅ N ₂ C ₃ , deg	118.37	117.05	119.03	118.00
C ₆ N ₁ C ₂ , deg	118.41	117.14	119.21	118.24
<i>E</i> _{LUMO} , eV	1.5284	1.5959	–2.1197	–2.1270
<i>E</i> _{HOMO} , eV	–5.1543	–5.2702	–3.5050	–3.5592
ΔE , eV	6.6827	6.8661	1.3853	1.4321

Table 4. Molecular geometry, bond distances, bond angles, torsional angles, the calculated HOMO and LUMO energies and energy gaps of BCP in neutral state

Bond/angle name	HF/3-21G*	HF/6-31G(<i>d</i>)	B3LYP/3-21G*	B3LYP/6-31G(<i>d</i>)
C ₂ –C ₃ , Å	1.449	1.459	1.454	1.460
N ₁ –N ₂ , Å	2.746	2.712	2.744	2.717
C ₁ –N ₁ , Å	1.307	1.298	1.334	1.324
C ₂ –N ₁ , Å	1.345	1.343	1.360	1.353
C ₃ –N ₂ , Å	1.345	1.343	1.360	1.353
C ₄ –N ₂ , Å	1.307	1.298	1.334	1.324
N ₁ C ₂ C ₃ N ₂ , deg	0.000	0.000	0.000	0.000
C ₇ C ₁ N ₁ C ₂ , deg	0.334	0.637	0.712	0.917
C ₈ C ₄ N ₂ C ₃ , deg	–0.334	–0.637	–0.712	–0.917
C ₁ N ₁ C ₂ , deg	120.47	119.38	119.23	118.80
C ₃ N ₂ C ₄ , deg	120.47	119.38	119.23	118.80
C ₆ C ₁ N ₁ , deg	118.30	117.86	117.60	117.31
C ₅ C ₄ N ₂ , deg	118.30	117.86	117.60	117.31
<i>E</i> _{LUMO} , eV	2.4704	2.4351	–1.2947	–1.2849
<i>E</i> _{HOMO} , eV	–7.9804	–7.7556	–5.9385	–5.7869
ΔE , eV	10.4509	10.1907	4.6438	4.5020

Table 5. Molecular geometry, bond distances, bond angles, torsional angles, the calculated HOMO and LUMO energies and energy gaps of IM2⁺ in neutral state

Bond/angle name	HF/3-21G*	HF/6-31G(<i>d</i>)	B3LYP/3-21G*	B3LYP/6-31G(<i>d</i>)
C ₂ –C ₃ , Å	0	1.458	1.451	1.457
N ₁ –N ₂ , Å	2.866	2.866	2.854	2.864
C ₁ –N ₁ , Å	1.305	1.297	1.333	1.324
C ₂ –N ₁ , Å	1.343	1.341	1.354	1.349
C ₃ –N ₂ , Å	1.386	1.390	1.395	1.398
C ₄ –N ₂ , Å	1.346	1.340	1.370	1.364
C ₆ –N ₂ , Å	1.474	1.464	1.481	1.469
C ₇ –C ₁ , Å	3.079	3.090	3.091	3.115
N ₁ C ₂ C ₃ N ₂ , deg	–1.434	–12.163	–4.700	–9.929
C ₁ N ₁ C ₂ , deg	121.41	120.10	120.22	119.69
C ₃ N ₂ C ₄ , deg	121.61	121.20	121.65	121.44
C ₆ N ₂ C ₃ , deg	123.24	122.87	123.18	123.26
C ₈ C ₁ N ₁ , deg	118.91	118.46	118.14	117.84
C ₅ C ₄ N ₂ , deg	120.42	121.06	119.85	120.51
<i>E</i> _{LUMO} , eV	–2.1485	–2.0527	–5.5540	–5.4715
<i>E</i> _{HOMO} , eV	–9.6958	–9.4498	–7.6092	–7.4615
ΔE , eV	7.5472	7.3970	2.0552	1.9899

Table 6. Molecular geometry, bond distances, bond angles, torsional angles, the calculated HOMO and LUMO energies and energy gaps of D2 in neutral state

Bond/angle name	HF/3-21G*	HF/6-31G(d)	B3LYP/3-21G*	B3LYP/6-31G(d)
C ₂ –C ₃ , Å	1.324	1.328	1.357	1.361
C ₁ –N ₁ , Å	1.402	1.401	1.416	1.407
C ₂ –N ₁ , Å	1.423	1.417	1.436	1.430
C ₃ –N ₂ , Å	1.423	1.416	1.437	1.429
C ₄ –N ₂ , Å	1.402	1.401	1.407	1.405
C ₆ –N ₂ , Å	1.417	1.416	1.433	1.427
C ₆ –C ₇ , Å	1.405	1.402	1.416	1.410
C ₇ –N ₁ , Å	1.417	1.415	1.423	1.423
N ₁ C ₂ C ₃ N ₂ , deg	0.028	0.036	0.054	–0.082
N ₁ C ₇ C ₆ N ₂ , deg	0.293	0.432	0.513	0.837
C ₅ C ₄ N ₂ C ₆ , deg	35.300	–50.815	–36.703	–50.714
C ₈ C ₁ N ₁ C ₇ , deg	–35.245	–50.993	44.726	50.476
C ₁ N ₁ C ₂ , deg	118.02	117.76	117.95	117.81
C ₃ N ₂ C ₄ , deg	118.05	117.79	118.12	117.87
C ₈ C ₁ N ₂ , deg	118.91	117.67	118.42	117.92
C ₅ C ₄ N ₂ , deg	118.93	117.69	119.25	117.95
C ₂ N ₁ C ₇ , deg	115.37	113.15	115.61	113.94
C ₃ N ₂ C ₆ , deg	115.32	113.03	115.43	113.67
E _{LUMO} , eV	–1.6062	–1.6049	–2.0138	–2.0555
E _{HOMO} , eV	–5.2773	–5.5216	–3.5885	–3.7124
ΔE, eV	3.6710	3.9167	1.5747	1.6568

a neutral state [23]. The potential energy curve of this reaction is shown in Fig. 8 [22].

The hopping rates for the self-exchange reactions are given by Eq. (1), the semi-classical Marcus theory [24–28]

$$k_{\text{et}} = \left(\frac{4\pi^2}{h} \right) \Delta H_{\text{ab}}^2 (4\pi\lambda_{\pm} T)^{-\frac{1}{2}} e^{-\frac{\lambda_{\pm}}{4k_{\text{B}}T}}. \quad (1)$$

Here, λ_{\pm} are reorganization energies, ΔH_{ab} is the electronic coupling matrix element between donor and acceptor k_{B} is the Boltzman constant, and h is the Planck constant [24–28].

The reorganization energies are purely intrinsic properties of a single molecule and these energies consists of four terms corresponding to the structural relaxation energies λ_1 , λ_2 , λ_3 and λ_4 [22]. The sum of the relaxation energies λ_1 and λ_2 toward optimum geometries is the reorganization energy, λ_+ , for the hole transport process and the sum of the relaxation energies λ_3 and λ_4 toward optimum geometries is the reorganization energy, λ_- , for the electron transport process [22].

In the reorganization process λ_1 , the geometry changes from the optimized neutral state to the opti-

mized cationic state while the compound is in its cationic state [22]. In the λ_2 process, the geometry change is the reverse of λ_1 while the compound is in its neutral state [22]. Because these two processes involve the same geometrical change, it is reasonable to expect that λ_1 and λ_2 to be comparable. Accordingly, the reorganization energies λ_3 and λ_4 should be similar [22].

Table 7. Bond distances, bond angles and torsional angles of BPhen in cationic and anionic states, calculated by the DFT/B3LYP method with the 6-31G(d) basis set

Bond/angle name	Cationic	Anionic
C ₂ –C ₃ , Å	1.473	1.477
N ₁ –N ₂ , Å	2.710	2.710
C ₁ –N ₁ , Å	1.339	1.340
C ₂ –N ₁ , Å	1.330	1.340
C ₃ –N ₂ , Å	1.330	1.340
C ₄ –N ₂ , Å	1.339	1.340
N ₁ C ₂ C ₃ N ₂ , deg	0.000	0.000
C ₁ N ₁ C ₂ , deg	117.92	118.02
C ₃ N ₂ C ₄ , deg	117.92	118.02

Table 8. Bond distances, bond angles and torsional angles of BCP in cationic and anionic states, calculated by the DFT/B3LYP method with the 6-31G(*d*) basis set

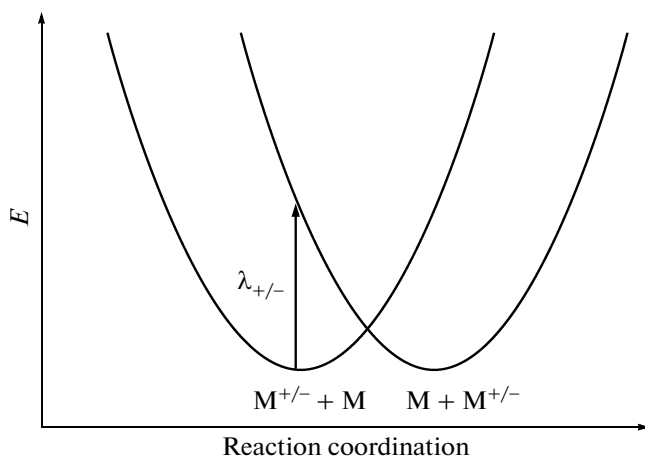
Bond/angle name	Cationic	Anionic
C ₂ –C ₃ , Å	1.478	1.478
N ₁ –N ₂ , Å	2.710	2.710
C ₁ –N ₁ , Å	1.350	1.345
C ₂ –N ₁ , Å	1.326	1.338
C ₃ –N ₂ , Å	1.326	1.338
C ₄ –N ₂ , Å	1.350	1.345
N ₁ C ₂ C ₃ N ₂ , deg	0.000	0.000
C ₇ C ₁ N ₁ C ₂ , deg	0.458	2.089
C ₈ C ₄ N ₂ C ₃ , deg	–0.458	–2.089
C ₁ N ₁ C ₂ , deg	118.90	118.91
C ₃ N ₂ C ₄ , deg	118.90	118.91
C ₆ C ₁ N ₁ , deg	115.86	116.04
C ₅ C ₄ N ₂ , deg	115.86	116.04

Table 11 shows reorganization energies λ_+ and λ_- of BPhen, BCP, D1 and D2 dimers plus their components λ_1 , λ_2 , λ_3 and λ_4 .

The values of λ_1 and λ_2 are nearly equal for the monomers and the D1 dimer but are quite different for the D2 dimer. The magnitudes of λ_3 and λ_4 for the monomers and the dimers are nearly equal.

The large difference between λ_1 and λ_2 for the D2 dimer can be explained by investigating geometrical changes in its anionic and cationic states.

When the D2 dimer state changes from neutral to cationic, its torsional angles C₅C₄N₂C₆, C₈C₁N₁C₇ and N₁C₂C₃N₂ change from -50.714° to -29.608° , 50.476° to 29.353° and -0.082° to 0.019° respectively and the C₄–N₂ and C₁–N₁ bonds are strengthened

**Fig. 8.** Schematic diagram in the transitions for the hole and electron transports [22].**Table 9.** Bond distances, bond angles and torsional angles of D1 in cationic and anionic states, calculated by the DFT/B3LYP method with the 6-31G(*d*) basis set

Bond/angle name	Cationic	Anionic
C ₂ –C ₃ , Å	1.382	1.395
C ₁ –N ₁ , Å	1.407	1.374
C ₂ –N ₁ , Å	1.434	1.405
C ₃ –N ₂ , Å	1.427	1.405
C ₄ –N ₂ , Å	1.394	1.374
C ₅ –N ₂ , Å	1.410	1.427
C ₅ –C ₆ , Å	1.434	1.407
C ₆ –N ₁ , Å	1.382	1.424
N ₁ C ₂ C ₃ N ₂ , deg	–0.479	0.124
N ₂ C ₅ C ₆ N ₁ , deg	3.798	0.961
C ₂ C ₃ N ₂ C ₅ , deg	25.251	16.655
C ₁ N ₁ C ₂ , deg	117.71	119.16
C ₃ N ₂ C ₄ , deg	118.42	119.19
C ₅ N ₂ C ₃ , deg	118.87	119.21
C ₆ N ₁ C ₂ , deg	119.82	119.27

(Table 6 and 10). This bond strengthening impedes the relaxation of the C₅C₄N₂C₆ and C₈C₁N₁C₇ torsional angles. However, the C₂–C₃ bond is weakened, aiding the relaxation of N₁C₂C₃N₂ torsional angle. These torsional angle changes are highly significant and are dominant factors in determining λ_1 and λ_2 energies.

In the λ_2 process, the energy required to reorganize C₅C₄N₂C₆ and C₈C₁N₁C₇ angles is greater than the energy required in the λ_1 process.

When the D2 dimer state changes from neutral to anionic, its torsional angles C₅C₄N₂C₆, C₈C₁N₁C₇ and N₁C₂C₃N₂ change from -50.714° to -55.442° , 50.476° to 50.004° and -0.082° to 0.0708° respectively (Table 6 and 10). These angle changes are negligible.

There is a minor difference between λ_1 and λ_2 in the D1 dimer. For the D1 dimer in its cationic state, the C₃–N₂ bond is weakened and this bond weakening will facilitate the C₂C₃N₂C₅ angle to change from 26.715° to 16.655° .

There is a small difference between λ_3 and λ_4 in the BCP monomer. For the BCP monomer in its anionic state, the C₁–N₁ bond is weakened and this bond weakening will facilitate the dihedral angle C₇C₁N₁C₂ to change from 0.917° to 2.089° .

If one assumes T to be 298.15 K and neglects the small differences in ΔH_{ab} for electrons and holes, then a modified form of Eq. (1) will be [17]:

$$k_{et} = (\lambda_+/\lambda_-)^{1/2} \exp[(\lambda_+ - \lambda_-)/(4k_B T)]. \quad (2)$$

Table 10. Bond distances, bond angles, and torsional angles of D2 in cationic and anionic states, calculated by the DFT/B3LYP method with the 6-31G(*d*) basis set

Bond/angle name	Cationic	Anionic
C ₂ –C ₃ , Å	1.375	1.393
C ₁ –N ₁ , Å	1.421	1.388
C ₂ –N ₁ , Å	1.435	1.410
C ₃ –N ₂ , Å	1.434	1.410
C ₄ –N ₂ , Å	1.418	1.387
C ₆ –N ₂ , Å	1.413	1.430
C ₆ –C ₇ , Å	1.420	1.406
C ₇ –N ₁ , Å	1.399	1.428
N ₁ C ₂ C ₃ N ₂ , deg	–0.708	0.019
N ₁ C ₇ C ₆ N ₂ , deg	2.227	0.828
C ₅ C ₄ N ₂ C ₆ , deg	–55.442	–29.608
C ₈ C ₁ N ₁ C ₇ , deg	50.004	29.353
C ₁ N ₁ C ₂ , deg	117.65	119.34
C ₃ N ₂ C ₄ , deg	117.74	119.39
C ₈ C ₁ N ₁ , deg	117.57	120.51
C ₅ C ₄ N ₂ , deg	117.46	120.54
C ₂ N ₁ C ₇ , deg	115.52	115.78
C ₃ N ₂ C ₆ , deg	114.56	115.63

Table 11 shows the relative hopping rates, ($k_{\text{et-}}/k_{\text{et+}}$), of electrons over holes calculated using the Eq. (2).

The relative hopping rates data show that the dimerization reaction of the BCP monomer causes degradation of the electron transporting BCP under hole transporting conditions. The reason for this degradation is because the relative hopping rates of the D2 dimer is larger than its monomer and increases as the amount of phenyl rings in the D2 dimer increases.

The relative hopping rates of the D1 dimer is smaller than its monomer and decreases as the amount of phenyl rings in the D1 dimer increases. The calculations show that the order of $k_{\text{et-}}/k_{\text{et+}}$ is D2 > BPhen > BCP > D1 and thus the presence of methyl groups in

the D1 and D2 dimers seems to make the $k_{\text{et-}}/k_{\text{et+}}$ value to be lowest and highest, respectively.

Table 11 shows the magnitude of λ_- is larger than λ_+ for all compounds except for the D2 dimer. The value of λ_- for the BCP monomer is larger than λ_- of the BPhen monomer and λ_- of the D1 dimer is larger than the λ_- of the BPhen monomer but the value of the D2 dimer is smaller than the λ_- of the BCP monomer. These results show that the methyl groups and the phenyl rings act as sterically hindering groups.

3.3. Ionization Potential and Electron Affinity

Ionization potentials (IP), Electron Affinities (EA), Re-organization energies λ_+ and λ_- , together with its Components, λ_1 , λ_2 , λ_3 and λ_4 , calculated by the DFT/B3LYP method with the 6-31G(*d*) basis set are given in Table 11.

The higher IP of the BPhen monomer compared to that of the BCP monomer is due to the electron donating effect of the methyl groups in the BCP monomer.

The magnitudes of IP for the dimers are smaller than the IP values for their monomers. The values of IP (7.37 and 7.15 eV) and ΔE (4.50 and 4.50 eV) for the BPhen and the BCP monomers are large compared with values of IP (4.74 and 4.87 eV) and ΔE (1.43 and 1.65 eV) for the D1 and D2 dimers.

The slightly smaller EA of the BCP monomer than that of the BPhen monomer shows that the methyl groups are effective in lowering EA. The magnitudes of EA and HOMO energies for the dimers are larger than the corresponding magnitudes in their monomers but the magnitudes of IP and LUMO for the monomers are larger than the corresponding magnitudes in their dimers.

3.4. Thermodynamic Properties

Table 12 shows Zero-point vibrational energy, entropy, sum of electronic and thermal free energy, sum of electronic and thermal enthalpy and total energy were obtained for the monomers, positively charged intermediates and the dimers at the

Table 11. Relative hopping rates of electron over hole, Ionization potentials (IP) Electron Affinities (EA), Reorganization energies λ_+ and λ_- , together with its Components, λ_1 , λ_2 , λ_3 , and λ_4 , calculated by the DFT/B3LYP method with the 6-31G(*d*) basis set

	$k_{\text{et-}}/k_{\text{et+}}^a$	IP, eV	EA, eV	λ_1 , eV	λ_2 , eV	λ_+ , eV	λ_3 , eV	λ_4 , eV	λ_- , eV
BPhen	0.5267	7.3693	0.1081	0.1615	0.1617	0.3232	0.1836	0.1970	0.3806
BCP	0.4904	7.1567	–0.0165	0.1798	0.1874	0.3672	0.1953	0.2368	0.4321
D1	0.4410	4.7452	1.0511	0.1480	0.1702	0.3182	0.2005	0.1912	0.3917
D2	6.1009	4.8778	0.9764	0.2170	0.3148	0.5318	0.1808	0.1841	0.3649

^a Calculated for 298.15 K, neglecting differences in ΔH_{ab} .

Table 12

	HF/3-21G*	HF/6-31G(<i>d</i>)	B3LYP/3-21G*	B3LYP/6-31G(<i>d</i>)
Bphen	225.9683 ^a	223.9130 ^a	210.4390 ^a	208.9489 ^a
	139.641 ^b	139.247 ^b	141.657 ^b	142.845 ^b
	-1021.0304 ^c	-1026.7817 ^c	-1027.7309 ^c	-1033.4338 ^c
	-1020.9640 ^d	-1026.7156 ^d	-1027.6636 ^d	-1033.3660 ^d
	-1021.3421 ^e	-1027.0908 ^e	-1028.0182 ^e	-1033.7186 ^e
IM1 ⁺	446.9160 ^a	442.8486 ^a	416.0583 ^a	413.281 ^a
	234.933 ^b	234.201 ^b	240.202 ^b	241.464 ^b
	-2041.2902 ^c	-2052.6674 ^c	-2054.5404 ^c	-2065.9331 ^c
	-2041.1786 ^d	-2052.7787 ^d	-2054.6545 ^d	-2066.04792 ^d
	-2041.9267 ^e	-2053.4098 ^e	-2055.2420 ^e	-2066.6310 ^e
D1	438.8049 ^a	434.8987 ^a	407.8433 ^a	405.096 ^a
	226.031 ^b	226.822 ^b	234.050 ^b	234.283 ^b
	-2054.2251 ^c	-2052.2115 ^c	-2054.2251 ^c	-2065.4985 ^c
	-2040.8409 ^d	-2052.3192 ^d	-2054.1139 ^d	-2065.6098 ^d
	-2041.4680 ^e	-2052.9406 ^e	-2054.8018 ^e	-2066.1827 ^e
BCP	262.4633 ^a	260.3569 ^a	244.9912 ^a	243.2672 ^a
	156.704 ^b	157.184 ^b	159.531 ^b	161.633 ^b
	1098.6277 ^c	-1104.8106 ^c	-1105.8971 ^c	-1111.9498 ^c
	-1098.5532 ^d	-1104.7360 ^d	-1105.8213 ^d	-1112.0266 ^d
	-1098.9931 ^e	-1105.1729 ^e	-1106.2347 ^e	-1112.3610 ^e
IM2 ⁺	520.3266 ^a	516.089 ^a	485.4428 ^a	482.1507 ^a
	266.428 ^b	266.833 ^b	273.414 ^b	277.313 ^b
	-2196.4857 ^c	-2208.7067 ^c	-2210.9895 ^c	-2223.2332 ^c
	-2196.3591 ^d	-2208.8335 ^d	-2210.8596 ^d	-2223.1014 ^d
	-2197.2312 ^e	-2209.5728 ^e	-2211.6791 ^e	-2223.9164 ^e
D2	512.0578 ^a	508.5772 ^a	477.0269 ^a	474.3969 ^a
	258.224 ^b	258.916 ^b	269.924 ^b	267.711 ^b
	-2195.8941 ^c	-2208.3593 ^c	-2210.5408 ^c	-2222.7780 ^c
	-2196.0168 ^d	-2208.2363 ^d	-2210.4126 ^d	-2222.6508 ^d
	-2196.7524 ^e	-2209.0897 ^e	-2211.2181 ^e	-2223.4525 ^e

^a Zero-point vibrational energy (Kcal mol⁻¹), ^b entropy (cal mol⁻¹ K⁻¹), ^c sum of electronic and thermal free energy (hartree), ^d sum of electronic and thermal enthalpy (hartree) and ^e total energy (hartree).

HF/3-21G*, HF/6-31G(*d*), B3LYP/3-21G* and B3LYP/6-31G(*d*) levels of ab initio method. The thermodynamic properties, calculated at the same level of ab initio method with different basis sets, are only slightly different. Total energies and triplet energies of excited positively charged monomers are given in Table 13. Table 14 exhibits relative total energies (relative to the reactants) for reactions (1) and (2) at the levels of HF/3-21G*, HF/6-31G(*d*), B3LYP/3-21G* and B3LYP/6-31G(*d*). The energy variation in the reactions (1) and (2), calculated by the B3LYP/6-31G(*d*) method, is depicted in Fig. 9. This figure shows that total energies of IM1, IM2, P1 and P2 are 22.20,

26.66, 303.50 and 313.77 kcal mol⁻¹ and they are higher than the total energies of their reactants R1 and R2. It is clear that the different methods predict different values and since there are no experimental data are available for all compounds to be compared with our theoretical results (Tables 12–14), hence, we consider the present results with B3LYP/6-31G(*d*) are more accurate, DFT methods such as B3LYP/6-31G(*d*) are often times considered to be standard model chemistry for many applications. We hope that our present work will stimulate some more works on these materials. Future experimental works will testify our calculated results.

Table 13. Total energies and excitation energies for excited positively charged monomers (BPhen^{*+} and BCP^{*+}) at td-triplets

	HF/3-21G*	HF/6-31G(d)	B3LYP/3-21G*	B3LYP/6-31G(d)
BPhen ^{*+}	-1021.0704 ^a	-1026.8446 ^a	1027.7371 ^a	-1033.4477 ^a
	-2.2282 ^b	-2.2783 ^b	0.0999 ^b	0.3903 ^b
BCP ^{*+}	-1098.7269 ^a	-1104.9332 ^a	-1105.9612 ^a	-1112.0979 ^a
	-2.3957 ^b	-2.3076 ^b	0.0916 ^b	0.4114 ^b

^a Total energies (Hartree) and ^b excitation energies (eV).

Table 14. Relative total energies (relative to the reactants) for reactions (1) and (2), at the levels of HF/3-21G*, HF/6-31G(d), B3LYP/3-21G* and B3LYP/6-31G(d), in eV

	HF/3-21G*	HF/6-31G(d)	B3LYP/3-21G*	B3LYP/6-31G(d)
R1	0.00	0.00	0.00	0.00
IM1	-0.2827	0.7452	0.4356	0.9627
P1	12.1989	13.5126	12.4139	13.1609
R2	0.00	0.00	0.00	0.00
IM2	-0.2010	0.9542	0.5303	1.1564
P2	12.8384	14.0998	13.0745	13.7796

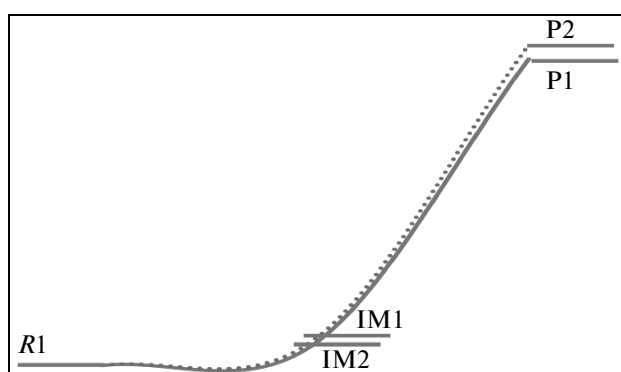
3.5. Conclusion

This theoretical work has helped us to gain further insight into the dimerization reactions of the BPhen and the BCP. The optimization of geometry calculations depend on the initial geometry around nitrogen atoms in the BPhen and the BCP reactants and their products. The variations in the geometries can be seen from HOMO and LUMO energies. The BPhen and the BCP dimers have the lowest LUMO energy levels in comparison with the LUMO energy levels for the BPhen and the BCP monomers and this is due to the dimers expanded *p*-electron system.

The reorganization energy λ_- of BPhen, BCP, BPhen intermediate, BCP intermediate and BPhen dimer is larger than λ_+ , except for the BCP dimer. The magnitudes of HOMO and EA for the dimers are

larger than the corresponding values in their monomers. The magnitudes of LUMO, IP and energy gap for the monomers are larger than the corresponding values in their dimers. The magnitudes of λ_- for the BPhen monomer is smaller than its λ_- for the BPhen dimer and λ_- of the BCP monomer is larger than λ_- of the BCP dimer. In contrast, the magnitude of λ_+ of the BPhen monomer is larger than its λ_+ value of the BPhen dimer and λ_+ value of the BCP monomer is smaller than its value in the BCP dimer. The relative hopping rates of electrons over holes for the BCP dimer is larger than its rate for its monomer, but for the BPhen dimer is smaller than its monomer.

Our results show that not only the methyl group, but also the phenyl rings act as sterically hindering groups and the generation of the BCP dimer is less pronounced than the BPhen dimer. The electron transport rate of the BPhen monomer is slightly higher than that of the BCP monomer. The electron transport rate of the BPhen dimer is much higher than that of the BCP monomer and the electron transport rate of the BCP dimer is lower than that of the BPhen monomer.

**Fig. 9.** Schematic energy profile of the reactions (P1) and (P2) calculated by B3LYP/6-31G method.

ACKNOWLEDGMENTS

This work was supported by Graduate University of Advanced Technology and Research Center for Science, High Technology & Environmental Science. This article was peer reviewed by Nasser Mirzai-Baghini, Imperial College Reactor Centre, Silwood Park, Ascot, Berkshire, United Kingdom.

REFERENCES

1. A. R. Duggal, J. J. Shiang, C. M. Heller, and D. F. Foust, *Appl. Phys. Lett.* **80**, 3470 (2002).
2. S. R. Forrest, *Nature (London)* **428**, 911 (2004).
3. Y. Shirota, K. Okumoto, and H. Inada, *Synth. Met.* **387**, 111 (2000).
4. B. C. Maiti, S. Z. Wang, C. P. Cheng, D. J. Huang, and C.-I. Chao, *Chin. Chem. Soc.* **48**, 1059 (2001).
5. S. Sebastian S. C. Cathrin, W. Karsten, K. Dirk, and L. Karl, *Org. Electron.* **8**, 709 (2007).
6. H. Aziz and Z. D. Popovic, *Chem. Mater.* **16**, 4522 (2004).
7. H. Aziz, Z. D. Popovic, N. Hu, A. Hor, and G. Xu, *Science* **283**, 1900 (1999).
8. Z. D. Popovic, H. Aziz, N.-X. Hu, A. Ioannidis, and P. N. M. dos Anjos, *J. Appl. Phys.* **89**, 4673 (2001).
9. F. Steuber, J. Staudigel, M. Stössel, J. Simmerer, A. Winnacker, H. Spreitzer, F. Weissörtel, and J. Salbeck, *Adv. Mater.* **12**, 130 (2000).
10. M. Müller-Wiegand, G. Georgiev, E. Oesterschulze, T. Fuhrmann, and J. Salbeck, *Appl. Phys. Lett.* **81**, 4940 (2002).
11. S. T. Lee, Z. Q. Gao, and L. S. Hung, *Appl. Phys. Lett.* **75**, 1404 (1999).
12. M. Scharnberg, J. Hu, J. Kanzow, K. Ratzke, R. Adlung, F. Faupel, C. Pannemann, U. Hilleringmann, S. Meyer, and J. Pflaum, *Appl. Phys. Lett.* **86**, 024104 (2005).
13. W. Song, Z. Li, S. K. So, Y. Qiu, Y. Zhu, and L. Cao, *Surf. Interface Anal.* **32**, 102 (2001).
14. M. J. Frisch, G. W. Trucks, H. B. Schlegel, et al., *Gaussian 03* (Gaussian Inc., Pittsburgh PA, 2003).
15. R. E. Stratmann, G. E. Scuseria, and M. J. Frisch, *J. Chem. Phys.* **109**, 8218 (1998).
16. S. Gao, *Comput. Phys. Commun.* **153**, 190 (2003).
17. K. Schwarz, *J. Solid State Chem.* **176**, 319 (2003).
18. B. B. Karki, L. Stixrude, S. J. Clark, M. C. Warren, G. J. Ackland, and J. Crain, *Am. Mineralog.* **82**, 635 (1997).
19. Li Zong-You, T. Shin-Rong, Ch. Yu-Chiang, et al., *Synth. Met.* **161**, 426 (2011).
20. A. Kahn, N. Koch, and W. Gao, *J. Polym. Sci. B: Polym. Phys.* **41**, 2529 (2003).
21. H. Bassler, *Phys. Status Solidi B* **175**, 15 (1993).
22. B. C. Lin, C. P. Cheng, and Z. P. M. Lao, *J. Phys. Chem. A* **107**, 5241 (2003).
23. H. P. Young, Y.-H. Kim, S. K. Kwon, I. S. Koo, and K. Yang, *Bull. Korean Chem. Soc.* **31**, 1649 (2010).
24. R. A. Marcus, *J. Chem. Phys.* **24**, 966 (1956).
25. R. A. Marcus, *J. Chem. Phys.* **43**, 679 (1965).
26. M. D. Newton and N. Sutin, *Ann. Rev. Phys. Chem.* **35**, 437 (1984).
27. R. A. Marcus and N. Sutin, *Biochim. Biophys. Acta* **811**, 265 (1985).
28. P. F. Barbara, T. J. Meyer, and M. A. Ratner, *J. Phys. Chem.* **100**, 13148 (1996).

SPELL: OK